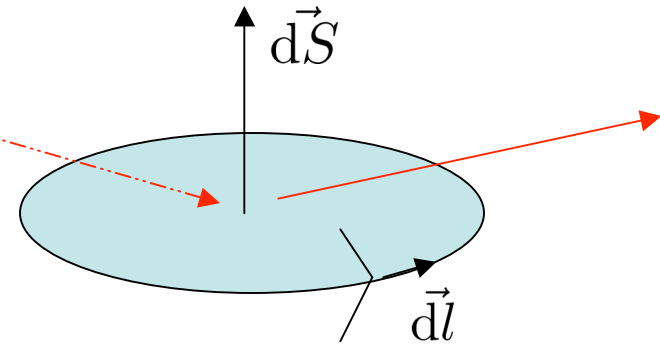


# Lecture 10: Curl

- Consider any vector field  $\vec{A}$
- And a loop in one plane with vector area  $d\vec{S}$
- If  $\vec{A}$  is **conservative** then

$$\oint \vec{A} \cdot d\vec{l} = 0 \text{ and } \vec{A} = \vec{\nabla} \phi$$



- But in general  $\oint \vec{A} \cdot d\vec{l} \neq 0$
- And we introduce a new (**pseudo**) vector
- Of magnitude:
- Which depends on the orientation of the loop
- The direction normal to the direction of the plane (of area  $=|d\vec{S}|$ ) which maximises the line integral

$$\text{curl} \vec{A} = \lim_{\text{area} \rightarrow 0} \frac{\oint \vec{A} \cdot d\vec{l}}{\text{area}}$$

Can evaluate 3 components by taking areas with normals in xyz directions

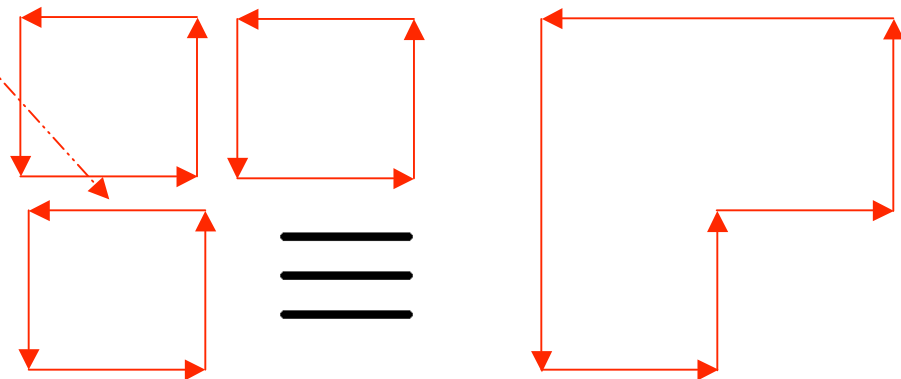
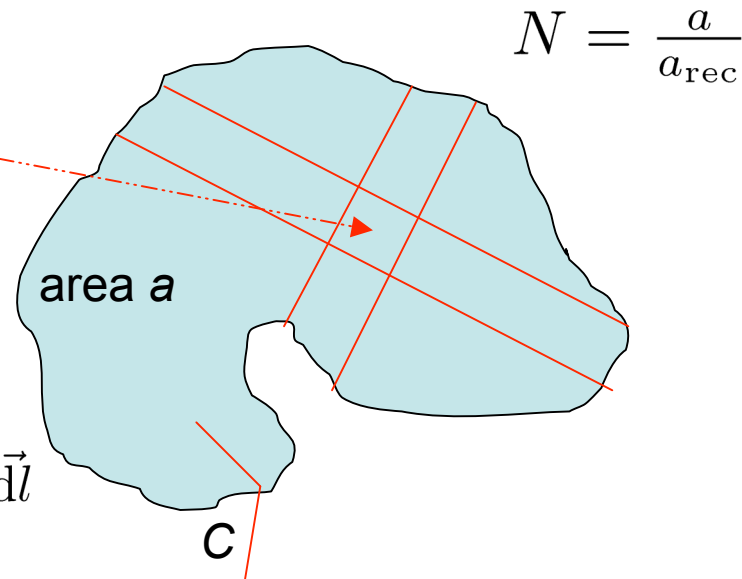
...is the direction of  $\text{curl} \vec{A}$  (will prove soon)  $\oint \vec{A} \cdot d\vec{l} = I_{\text{max}}$

# Curl is independent of (small) loop shape

- Sub-divide loop into  $N$  tiny rectangles of area  $a_{\text{rec}}$
- Summing line integrals, the interior contributions cancel, yielding a line integral over boundary  $C$

$$\frac{1}{a} \oint_C \vec{A} \cdot d\vec{l} = \sum_N \frac{1}{a} \oint_{\text{rec}} \vec{A} \cdot d\vec{l} = N \frac{a_{\text{rec}}}{a} \frac{1}{a_{\text{rec}}} \oint \vec{A} \cdot d\vec{l}$$

- Thus  $\oint \vec{A} \cdot d\vec{l} / \text{area}$  is same for large loop as for each rectangle
- Hence,  $\text{curl } \vec{A}$  is independent of loop shape
- But loop does need itself to be vanishingly small

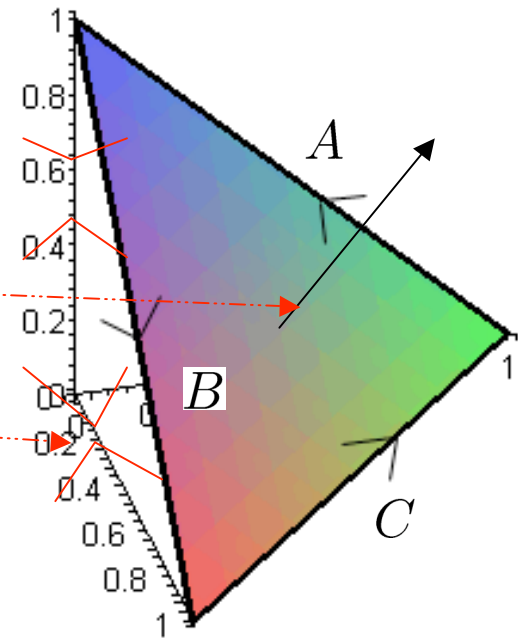


# Curl is a Vector

- Imagine a loop of arbitrary direction
- Make it out of three loops with normals in the  $\vec{i}, \vec{j}, \vec{k}$  directions

$$d\vec{S} = (dS_x, dS_y, dS_z)$$

- Because interior loops cancel, we get



$$\oint_{ABC} \vec{A} \cdot d\vec{l} = \oint_{A(-\vec{k})(\vec{j})} + \oint_{B(-\vec{i})(\vec{k})} + \oint_{C(-\vec{j})(\vec{i})} = C_x dS_x + C_y dS_y + C_z dS_z$$

$$\text{where } \text{curl} \vec{A} = (C_x, C_y, C_z)$$

$$\Rightarrow \oint \vec{A} \cdot d\vec{l} \equiv (\text{curl} \vec{A}) \cdot d\vec{S}$$

- Explains why line integral =  $I_{\text{max}} \cos \theta$

where  $\theta$  is angle between loop normal and direction of  $\text{curl} \vec{A}$

# In Cartesian Coordinates

- Consider tiny rectangle in  $yz$  plane

$$\oint \vec{A} \cdot d\vec{l} \text{ with } \vec{A} = (A_x, A_y, A_z) \\ d\vec{l} = (0, dy, dz)$$

- Contributions from sides 1 and 2

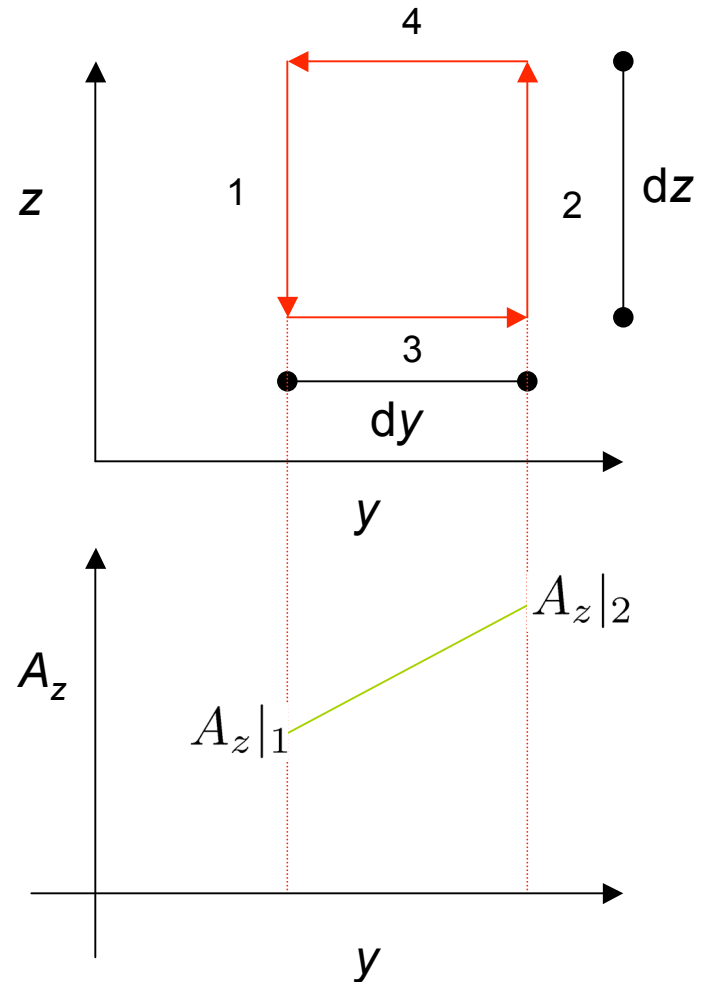
$$-(A_z|_1)dz + (A_z|_2)dz = dz [A_z|_2 - A_z|_1] = \\ dz \left[ dy \frac{\partial A_z}{\partial y} \right]$$

- Contributions from sides 3 and 4

$$-(A_y|_4)dy + (A_y|_3)dy = dy [A_y|_3 - A_y|_4] = \\ -dy \left[ dz \frac{\partial A_y}{\partial z} \right]$$

- Yielding in total

$$\oint_{yz} \vec{A} \cdot d\vec{l} = \left[ \frac{\partial A_z}{\partial y} - \frac{\partial A_y}{\partial z} \right] dy dz$$



- Repeating this exercise for a rectangular loop in the zx plane yields

$$\oint \vec{A} \cdot d\vec{l} = \left[ \frac{\partial A_x}{\partial z} - \frac{\partial A_z}{\partial x} \right] dx dz$$

- and in the xy plane yields

$$\oint \vec{A} \cdot d\vec{l} = \left[ \frac{\partial A_y}{\partial x} - \frac{\partial A_x}{\partial y} \right] dx dy$$

•

$$\oint \vec{A} \cdot d\vec{l} \equiv (\text{curl} \vec{A}) \cdot d\vec{S} = (C_x, C_y, C_z) \cdot d\vec{S}$$

$$\text{where } (C_x, C_y, C_z) = \left[ \frac{\partial A_z}{\partial y} - \frac{\partial A_y}{\partial z}, \frac{\partial A_x}{\partial z} - \frac{\partial A_z}{\partial x}, \frac{\partial A_y}{\partial x} - \frac{\partial A_x}{\partial y} \right]$$

# Key Equations

$$\oint \vec{A} \cdot d\vec{l} \equiv (\text{curl} \vec{A}) \cdot d\vec{S}$$

$$\text{curl} \vec{A} = \left[ \frac{\partial A_z}{\partial y} - \frac{\partial A_y}{\partial z}, \frac{\partial A_x}{\partial z} - \frac{\partial A_z}{\partial x}, \frac{\partial A_y}{\partial x} - \frac{\partial A_x}{\partial y} \right]$$

$$\text{curl} \vec{A} = \vec{\nabla} \times \vec{A} = \begin{vmatrix} \vec{i} & \vec{j} & \vec{k} \\ \frac{\partial}{\partial x} & \frac{\partial}{\partial y} & \frac{\partial}{\partial z} \\ A_x & A_y & A_z \end{vmatrix}$$

# Simple Examples

$$\begin{aligned}\vec{\nabla} \times \vec{r} &= \left( \frac{\partial}{\partial x}, \frac{\partial}{\partial y}, \frac{\partial}{\partial z} \right) \times (x, y, z) \\ &= \left( \frac{\partial z}{\partial y} - \frac{\partial y}{\partial z}, \frac{\partial x}{\partial z} - \frac{\partial z}{\partial x}, \frac{\partial y}{\partial x} - \frac{\partial x}{\partial y} \right) = (0, 0, 0)\end{aligned}$$

$$\begin{aligned}\vec{\nabla} \times (-y, x, z) &= \left( \frac{\partial z}{\partial y} - \frac{\partial x}{\partial z}, \frac{\partial(-y)}{\partial z} - \frac{\partial z}{\partial x}, \frac{\partial x}{\partial x} - \frac{\partial(-y)}{\partial y} \right) = (0, 0, 2)\end{aligned}$$

# Curl and Rotation

- Consider a **solid body** rotating with angular velocity  $\omega$  about the z axis

$$\vec{v} = \vec{\omega} \times \vec{r}$$

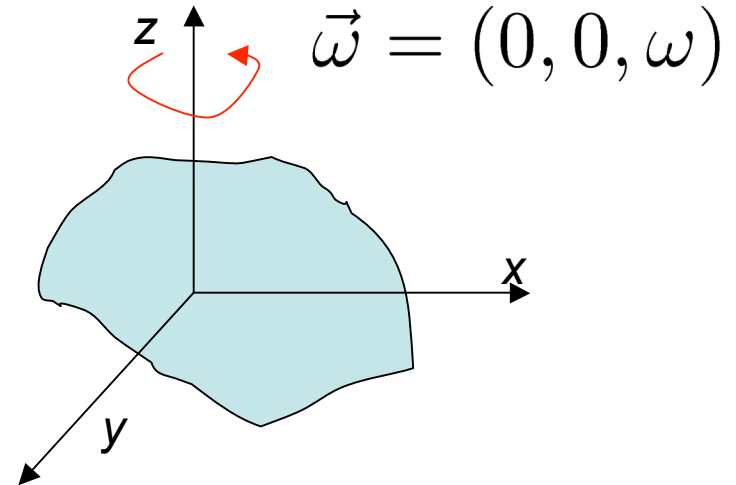
- describes the velocity vector field
- Evaluate curl  $\vec{v}$  from the definition
- x-component:

$\oint \vec{v} \cdot d\vec{l}$  in  $yz$  plane = 0 as  $\vec{v}$  is perpendicular to  $d\vec{l}$

- Similarly for y-component

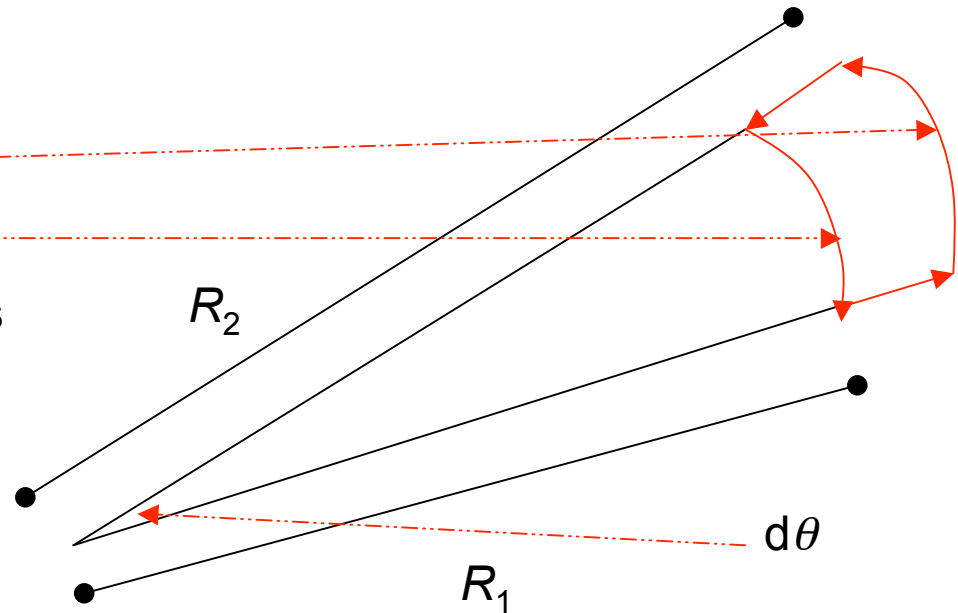
$\oint \vec{v} \cdot d\vec{l}$  in  $xz$  plane = 0 as  $\vec{v}$  is perpendicular to  $d\vec{l}$

- As only the z component is non-zero we see that  $\vec{\omega}$  and curl  $\vec{v}$  will be parallel





- z component
- In a rotating fluid, its easier to 'swim round' here than here
- i.e. **local evidence** that the fluid is rotating



$$[\text{curl } \vec{v}]_k \equiv \frac{(R_2\omega)R_2d\theta - (R_1\omega)R_1d\theta}{\frac{1}{2}(R_2^2 - R_1^2)d\theta} = 2\omega$$

- So magnitude of curl is proportional to any macroscopic rotation

If  $\text{curl } \vec{A} = 0 \Rightarrow$  Flow is said to be **irrotational**

- In a more general case, curl is a function of position